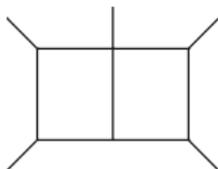


Physics at Hadron Colliders

Precision QCD at High Energies



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Department of Physics, University of Freiburg

SILAFEA
Antigua, Guatemala, November 2016



ALBERT-LUDWIGS-
UNIVERSITÄT FREIBURG



Alexander von Humboldt
Stiftung/Foundation

GREAT ACHIEVEMENTS @ LHC

Higgs, Complex signals, Complete SM, Beyond the SM?

SEARCHES FOR NEW PHYSICS

Bump hunting, Distribution excesses, Precision

THE QUEST FOR 1% QCD

The NNLO revolution, Th uncertainties, High multiplicities

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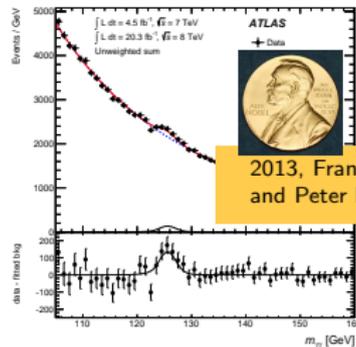
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A Complete Standard Model of Particle Physics

The SM is a quantum field theory that describes fundamental matter and their (strong and EW) interactions

- ▶ With the discovery of the Higgs Boson at the LHC in 2012, the last missing piece of the SM has been found
- ▶ We can now directly constraint all 19 parameters of the model

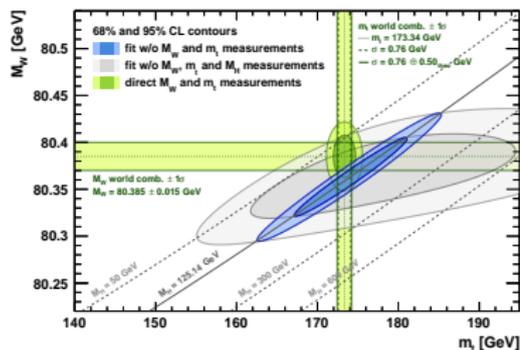
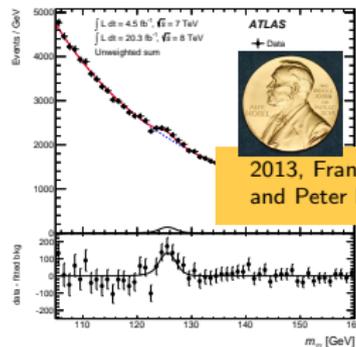


2013, François Englert and Peter Higgs

A Complete Standard Model of Particle Physics

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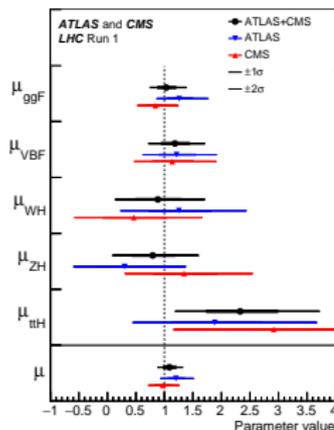
- ▶ With the discovery of the Higgs Boson at the LHC in 2012, the last missing piece of the SM has been found
- ▶ We can now directly constraint all 19 parameters of the model
- ▶ Global fits of observables can now be achieved, and theory/experiment comparisons can hint for problems with the SM
- ▶ The example figure shows a multidimensional fit by the Gfitter collaboration on the observables M_W , m_t and M_H



Higgs Phenomenology - LHC Run I

arXiv:1606.02266 [hep-ex]

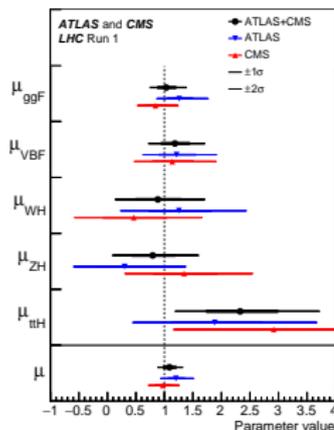
- ▶ Signal strengths for production mechanisms
- ▶ Includes *gluon fusion*, *VBF*, *Higgsstrahlung* and $t\bar{t}H$
- ▶ ATLAS, CMS and combined results shown



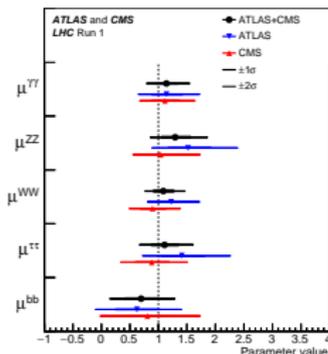
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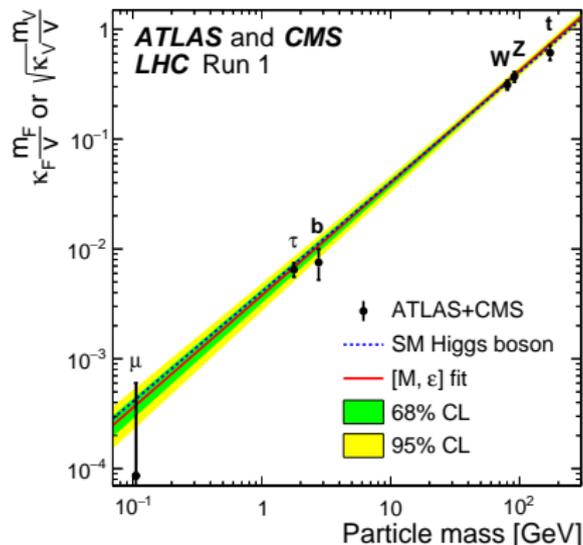
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- ▶ Includes *gluon fusion*, *VBF*, *Higgsstrahlung* and $t\bar{t}H$
- ▶ ATLAS, CMS and combined results shown



- ▶ Signal strengths for decay processes
- ▶ Shown are Higgs decaying into pairs of vector bosons and to fermion pairs
- ▶ Excellent overall agreement, though large uncertainties



Higgs Phenomenology - LHC Run I



The coupling strength of the Higgs boson to weak bosons (sqrt) and fermions as a function of the particle mass.

A *qualitative* compatibility to SM predictions is observed



A Tale of Troublesome Success

The Shortcomings of the SM

The Shortcomings of the SM

- ▶ Can we understand the **structure** of the SM symmetry group and its matter content?
- ▶ Why the **mass hierarchies** in the fermion sector and other peculiarities of its parameters?
- ▶ Why the **hierarchy** between the *electroweak* scale and the *Planck* scale?

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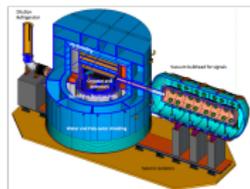
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Many Beyond the SM models have been proposed through the years to deal with some of these problems: *Supersymmetry, Extra dimensions, Composite Theories, Strings, Axions, Extra fields, etc.*

Searching for Answers: Dark Matter Experiments

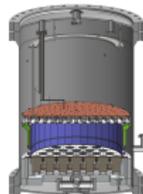
SuperCDMS: a cryogenic dark matter search experiment, located at SNOLAB, Ontario, Canada



CRESST: a cryogenic superconducting thermal dark matter experiment, located at Gran Sasso, Italy

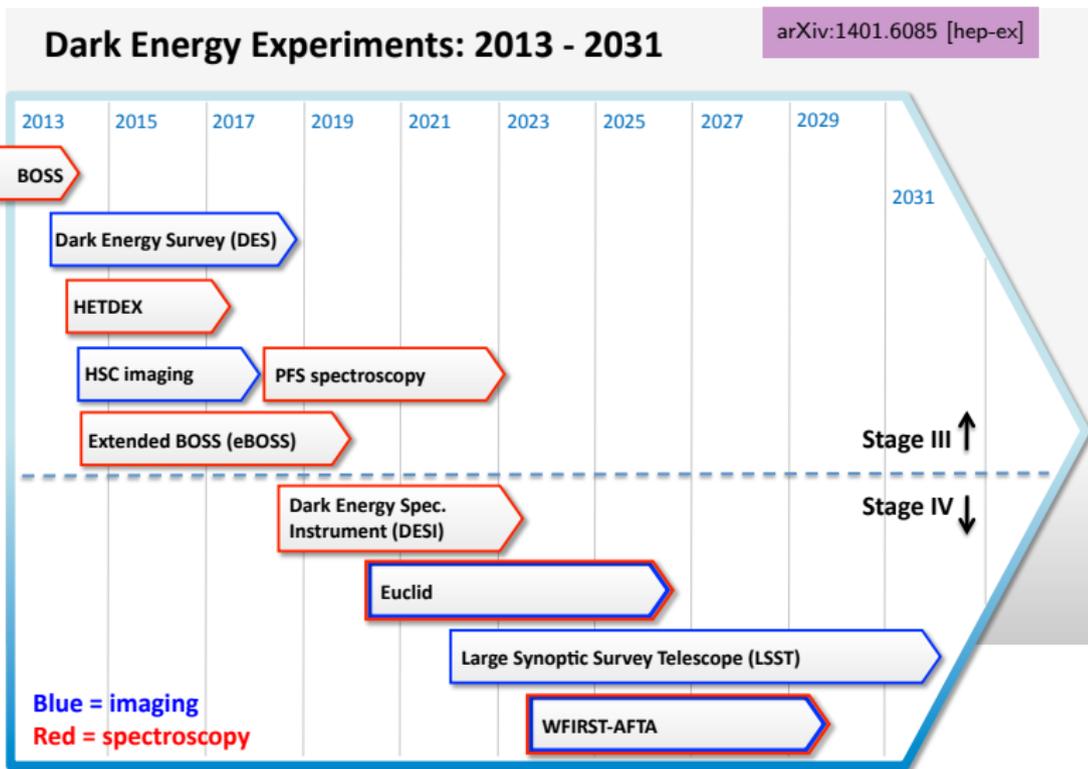


PandaX: a xenon based dark matter search experiment, located at CJPL, Sichuan, China



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Searching for Answers: Dark Energy Surveys



Searching for Answers with Collider Experiments

LHC: proton-proton collider up to a design energy of 14 TeV, CERN, Geneva



Searching for Answers with Collider Experiments

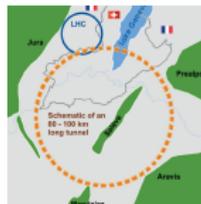
LHC: proton-proton collider up to a design energy of 14 TeV, CERN, Geneva



ILC: proposed Int. Linear Collider (0.5-1 TeV) with possible hosts Japan, Europe or the USA



FCC-CEPC: proposed future circular colliders to reach ~ 100 TeV in energy, possibly at CERN and China



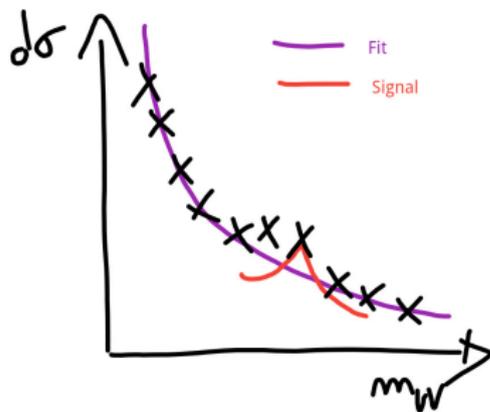
Clear New Signals

vs.

New Signal Excesses

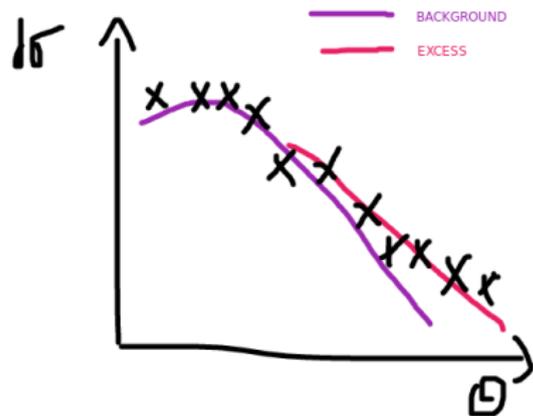
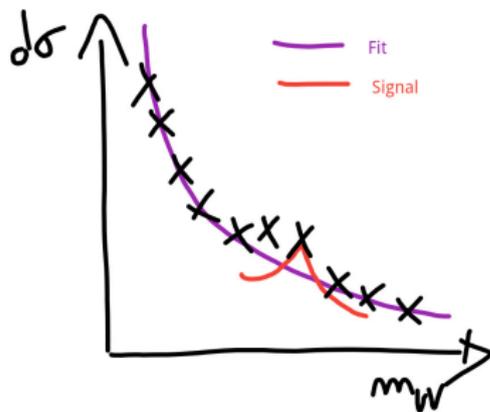
Bumps and Excesses at Colliders

When a heavy state is produced, that couples to SM particles, there is the possibility of a discovery by characterizing a *peak* on a related observable. Detection depends on the relative size of the signal and backgrounds



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When a heavy state is produced, that couples to SM particles, there is the possibility of a discovery by characterizing a *peak* on a related observable. Detection depends on the relative size of the signal and backgrounds



Unlike resonance signals, there are many BSM scenarios that enhance certain observables in a smooth way. In these cases a precise knowledge of the background is necessary

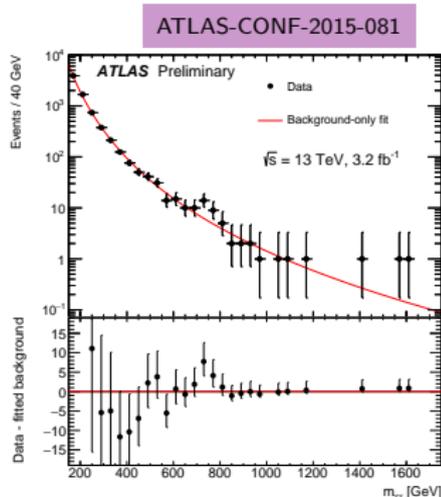
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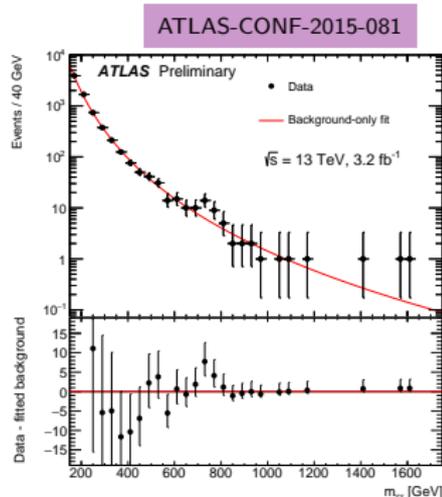
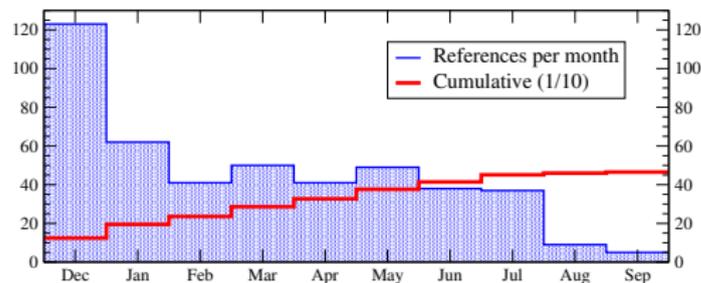
The Brief Story of a Diphoton Resonance

- ▶ On Dec. 15, 2015 the ATLAS and CMS collaboration reported on first results from Run II at $\sqrt{s} = 13$ TeV
- ▶ Both collaboration saw a curious excess of *diphoton* events at around $M_{\gamma\gamma} = 750$ GeV
- ▶ The statistical significance of the deviations was above 3 sigmas



The Brief Story of a Diphoton Resonance

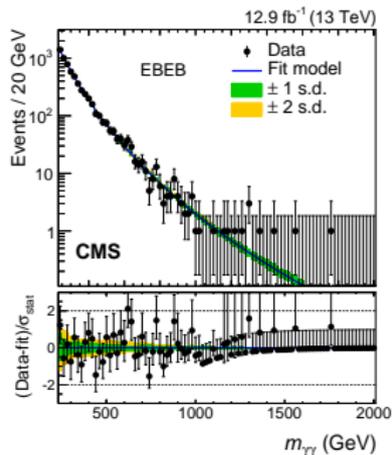
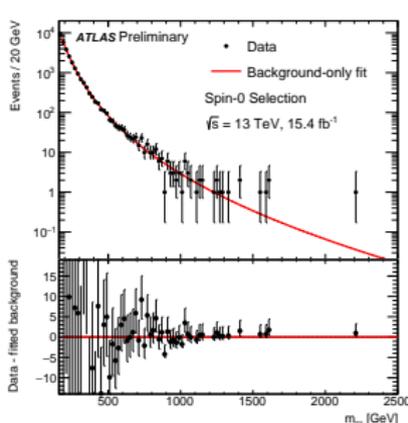
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- ▶ An avalanche of attention followed
- ▶ More than 400 articles explored in different ways the deviation

August Dismissal of a Fluctuation

But in Aug. 5, 2016 both collaborations revisited the measurement with more than 4 times the amount of data



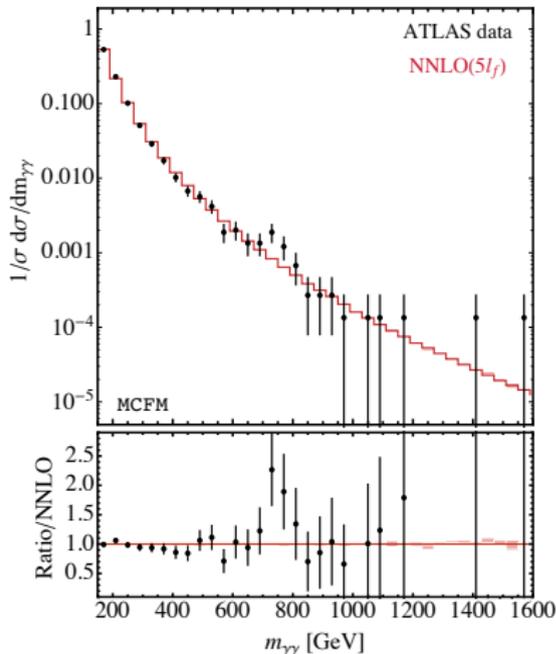
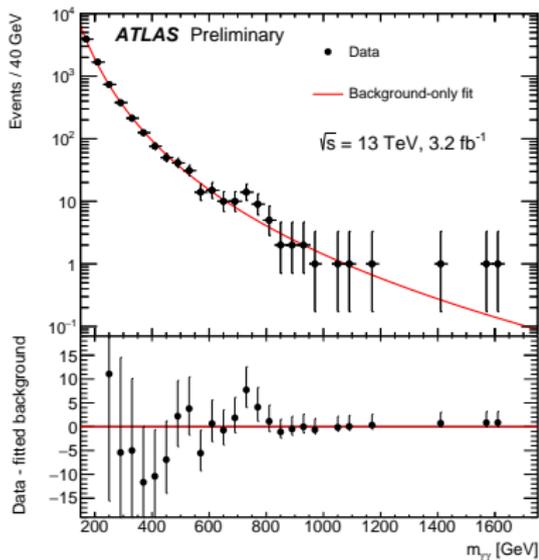
- ▶ The *evidence* of a resonance observed earlier was deemed a fluctuation of the background
- ▶ No matter how big data sets are, always in tails of distributions fake excesses can appear

Few Comments on the Experience

- ▶ Clear signatures are simplest to analyze, though dataset size important
- ▶ Eager community to find hints of BSM
- ▶ Precision Calculations are keen even for clear signatures, both to cross check experimental fits techniques and for finding $\mathcal{O}(1 - 5\%)$ effects

The $M_{\gamma\gamma}$ Spectrum at NNLO QCD

With a precise calculation by Campbell, Ellis, Li and Williams the experimental fits were validated



The $M_{\gamma\gamma}$ Spectrum at NNLO QCD

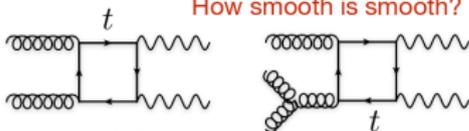
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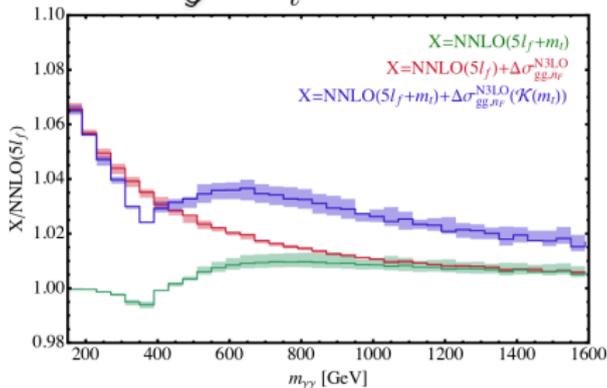
Predictions at high invariant masses.

As we all know, bump hunts in the diphoton system assume a smooth function which can be fitted to the data. Begging the question,

How smooth is smooth? :-)



Williams at Moriond 2016



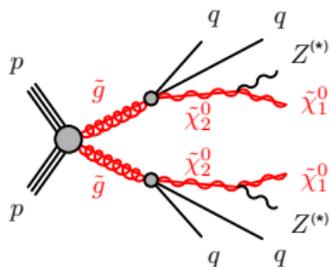
Searches at Hadron Colliders

Clear New Signals

vs.

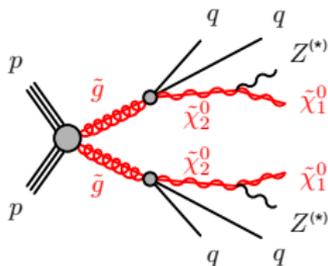
New Signal Excesses

Excesses in SUSY Searches

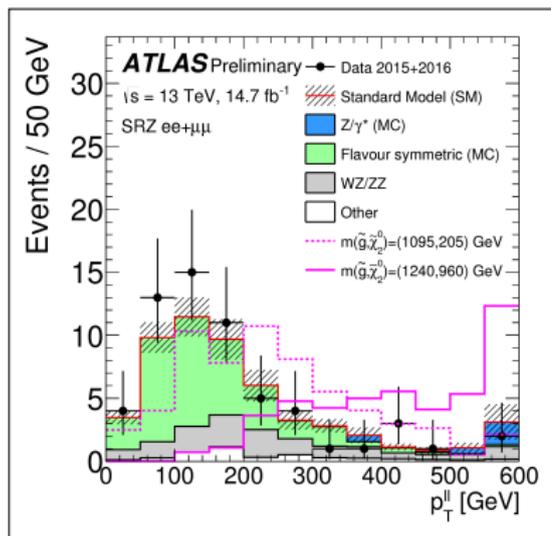


- ▶ As an example, in many SUSY models heavy colored particles are pair produced
- ▶ They produce long decay chains of jets and leptons
- ▶ In the end heavy neutral stable particles scape the detector, producing missing energy

Excesses in SUSY Searches



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p_T^{ll} is one of the employed observables, with possible NP modifying the associated distributions

Summary Plots for SUSY Exclusion Limits

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: August 2016

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13 \text{ TeV}$

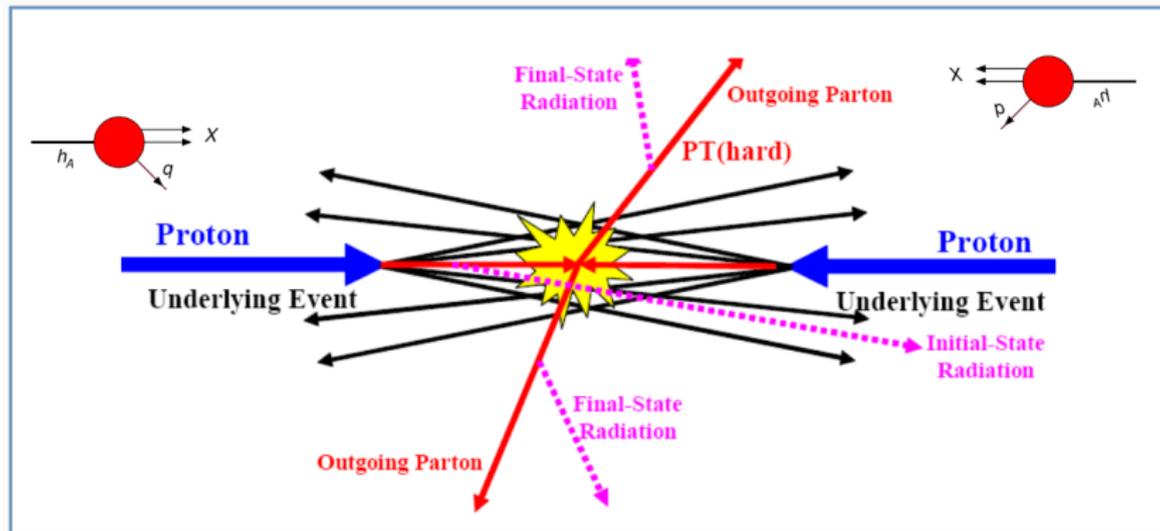
Model	e, μ, τ, γ Jets	E_{T}^{miss}	$\int \mathcal{L} d\mathcal{I} (\text{fb}^{-1})$	Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$	Reference	
Inclusive Searches	MSUGRA/CMSSM	0-3 e, μ, τ 2-10 jets/3b	Yes	20.3	4.3	1.85 TeV	$m(\tilde{g})=m(\tilde{t}_1)$	
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{g}$	0-2 jets	Yes	13.3	1	1.85 TeV	$m(\tilde{t}_1)=200 \text{ GeV}, m(\tilde{t}_2)=\text{max}(10, m(\tilde{t}_1)^2/\text{GeV})$	
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{t}_1$ (compressed)	mono-jet	Yes	3.2	1	608 GeV	$m(\tilde{g})=m(\tilde{t}_1)=5 \text{ GeV}$	
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{t}_2$	0-2 jets	Yes	13.3	1	1.85 TeV	$m(\tilde{t}_1)=0 \text{ GeV}$	
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{t}_1 + \tilde{W} \rightarrow q\tilde{t}_1 + \nu\tilde{W}$	0-2 jets	Yes	13.3	1	1.85 TeV	$m(\tilde{t}_1)=400 \text{ GeV}, m(\tilde{t}_2)=0.5 \text{ max}(10, m(\tilde{t}_1)^2/\text{GeV})$	
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{t}_1 + \tilde{W} \rightarrow q\tilde{t}_1 + \nu\tilde{W}$	3 e, μ, τ 4 jets	Yes	13.2	1	1.7 TeV	$m(\tilde{t}_1)=400 \text{ GeV}$	
	GMSB (bino NLSP)	2 e, μ (Z)	0-3 jets	Yes	13.2	1	1.6 TeV	$m(\tilde{t}_1)=500 \text{ GeV}$
	OGM (bino NLSP)	1-2 e, μ, τ 1-2 jets	Yes	3.2	1	2.0 TeV	$m(\text{NLSP})=0.1 \text{ mm}$	
	OGM (Higgsino-bino NLSP)	7	1 b	Yes	20.3	1	1.65 TeV	$m(\tilde{t}_1)=350 \text{ GeV}, m(\text{NLSP})=0.1 \text{ mm}, \mu=0$
	OGM (Higgsino-bino NLSP)	7	2 jets	Yes	13.3	1	1.8 TeV	$m(\tilde{t}_1)=600 \text{ GeV}, m(\text{NLSP})=0.1 \text{ mm}, \mu=0$
OGM (Higgsino NLSP)	2 e, μ (Z)	2 jets	Yes	20.3	1	1.85 TeV	$m(\text{NLSP})=400 \text{ GeV}$	
Gravitino LSP	0	mono-jet	Yes	20.3	1	900 GeV	$m(\tilde{g})=1.8 \times 10^{-4} \text{ eV}, m(\tilde{g})=m(\tilde{t}_1)=1.5 \text{ TeV}$	
3 γ gen. squarks & gluinos	$\tilde{g}, \tilde{q} \rightarrow q\tilde{g}$	0	3 b	Yes	14.8	1	1.89 TeV	$m(\tilde{t}_1)=0 \text{ GeV}$
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{t}_1$	0-1 e, μ, τ 3 b	Yes	14.8	1	1.89 TeV	$m(\tilde{t}_1)=0 \text{ GeV}$	
	$\tilde{g}, \tilde{q} \rightarrow q\tilde{t}_2$	0-1 e, μ, τ 3 b	Yes	20.1	1	1.37 TeV	$m(\tilde{t}_1)=300 \text{ GeV}$	
3 γ gen. squarks & gluinos (natural GMSB)	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0	2 b	Yes	3.2	1	840 GeV	$m(\tilde{t}_1)=100 \text{ GeV}$
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	2 e, μ (SS)	1 b	Yes	13.2	1	325-685 GeV	$m(\tilde{t}_1)=150 \text{ GeV}, m(\tilde{t}_2)=m(\tilde{t}_1)+100 \text{ GeV}$
	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0-2 e, μ, τ 1-2 b	Yes	4.7/13.3	1	200-720 GeV	$m(\tilde{t}_1)=2 \text{ max}(\tilde{t}_1), m(\tilde{t}_2)=0.5 \text{ GeV}$	
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$ or $\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0-2 e, μ, τ 0-2 jets 1-2 b	Yes	4.7/13.3	1	117-170 GeV	$m(\tilde{t}_1)=1 \text{ GeV}$	
	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0	mono-jet	Yes	3.2	1	90-323 GeV	$m(\tilde{t}_1)=m(\tilde{t}_2)=5 \text{ GeV}$
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	2 e, μ (Z)	1 b	Yes	20.3	1	150-600 GeV	$m(\tilde{t}_1)=100 \text{ GeV}$
	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	13.3	1	290-700 GeV	$m(\tilde{t}_1)=0 \text{ GeV}$
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2 + b$	1 e, μ, τ 6 jets + 2 b	Yes	20.3	1	320-620 GeV	$m(\tilde{t}_1)=0 \text{ GeV}$	
	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	2 e, μ, τ 0	Yes	20.3	1	90-935 GeV	$m(\tilde{t}_1)=0 \text{ GeV}$	
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	2 e, μ, τ 0	Yes	13.3	1	640 GeV	$m(\tilde{t}_1)=0 \text{ GeV}, m(\tilde{t}_2)=0.5 \text{ max}(\tilde{t}_1, m(\tilde{t}_2))$	
EW direct	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0	2 b	Yes	14.8	1	580 GeV	$m(\tilde{t}_1)=0 \text{ GeV}, m(\tilde{t}_2)=0.5 \text{ max}(\tilde{t}_1, m(\tilde{t}_2))$
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	3 e, μ, τ 0	Yes	13.3	1	1.0 TeV	$m(\tilde{t}_1)=m(\tilde{t}_2), m(\tilde{t}_3)=0, m(\tilde{t}_4)=0.5 \text{ max}(\tilde{t}_1, m(\tilde{t}_2))$	
	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	3 e, μ, τ 0	Yes	13.3	1	425 GeV	$m(\tilde{t}_1)=m(\tilde{t}_2), m(\tilde{t}_3)=0, \tilde{t}$ decoupled	
	$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	3 e, μ, τ 0-2 jets	Yes	20.3	1	270 GeV	$m(\tilde{t}_1)=m(\tilde{t}_2), m(\tilde{t}_3)=0, \tilde{t}$ decoupled	
	$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	3 e, μ, τ 0	Yes	20.3	1	935 GeV	$m(\tilde{t}_1)=m(\tilde{t}_2), m(\tilde{t}_3)=0, m(\tilde{t}_4)=0.5 \text{ max}(\tilde{t}_1, m(\tilde{t}_2))$	
	OGM (bino NLSP) weak prod.	1 e, μ, τ + γ	-	Yes	20.3	1	115-370 GeV	$\tau < 1 \text{ mm}$
	OGM (bino NLSP) weak prod.	2 γ	-	Yes	20.3	1	590 GeV	$\tau < 1 \text{ mm}$
	Direct \tilde{t}_1, \tilde{t}_2 prod., long-lived \tilde{t}_1	Disapp. trk	1 jet	Yes	3.2	1	270 GeV	$m(\tilde{t}_1)=m(\tilde{t}_2)=160 \text{ MeV}, m(\tilde{t}_1^2)=0.2 \text{ mm}$
	Direct \tilde{t}_1, \tilde{t}_2 prod., long-lived \tilde{t}_1	dE/dx trk	Yes	18.4	1	495 GeV	$m(\tilde{t}_1)=m(\tilde{t}_2)=160 \text{ MeV}, m(\tilde{t}_1^2)=15 \text{ ns}$	
	Stable, stopped \tilde{g} R-hadron	0-1 jets	Yes	27.9	1	850 GeV	$m(\tilde{t}_1)=100 \text{ GeV}, 10 \mu\text{s} < \tau < 1000 \mu\text{s}$	
Stable R-hadron	trk	-	3.2	1	1.58 TeV	1604-04520		
Metastable \tilde{g} R-hadron	dE/dx trk	-	3.2	1	1.57 TeV	10-damp-50		
GMSB, stable $\tilde{g}, \tilde{t}_1^0 \rightarrow \tilde{g}(\tilde{t}_1, \tilde{t}_2) + \nu(\mu, \tau)$	1-2 μ, τ	Yes	19.1	1	537 GeV	10-damp-50		
GMSB, $\tilde{t}_1^0 \rightarrow \tilde{g} + \gamma$, long-lived \tilde{t}_1^0	2 γ	Yes	20.3	1	440 GeV	1-c $\tilde{t}_1^0=3 \text{ ns}$, SPSe model		
GMSB, $\tilde{t}_1^0 \rightarrow \nu\tilde{g} + \nu\tilde{g}$	disapp. trk/spj/psj	-	3.2	1	1.0 TeV	7-c $\nu\tilde{t}_1^0=740 \text{ mm}, m(\tilde{t}_1)=1.3 \text{ TeV}$		
LFV $\tilde{g}, \tilde{q} \rightarrow q\tilde{g}\tilde{q}$, $\tilde{g}, \tilde{q} \rightarrow q\tilde{g}\tilde{q}$	pp ν jet	-	20.3	1	1.9 TeV	6-c $\nu\tilde{t}_1^0=480 \text{ mm}, m(\tilde{t}_1)=1.1 \text{ TeV}$		
Bilinear RPV CMSSM	2 e, μ (SS)	0-3 jets	Yes	20.3	1	1.45 TeV	$A_{11} < 0.11, A_{21} < 0.00000007$	
$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	4 e, μ, τ	Yes	13.3	1	1.14 TeV	$m(\tilde{t}_1)=4000 \text{ GeV}, A_{11} < 0.01, A_{21} < 0.00000007$		
$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	3 e, μ, τ + τ	Yes	20.3	1	450 GeV	$m(\tilde{t}_1)=0.2 \text{ max}(\tilde{t}_1), A_{11} < 0$		
$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0	4.5 large-R jets	Yes	14.8	1	1.08 TeV	$\text{BR}(\tilde{g} \rightarrow \tilde{g}\tilde{g}) = \text{BR}(\tilde{g} \rightarrow \tilde{g}\tilde{g}) = 0\%$	
$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	0	4.5 large-R jets	Yes	14.8	1	1.55 TeV	$m(\tilde{t}_1)=800 \text{ GeV}$	
$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	1 e, μ, τ 8-10 jets 0-4 b	Yes	14.8	1	1.75 TeV	$m(\tilde{t}_1)=700 \text{ GeV}$		
$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	1 e, μ, τ 8-10 jets 0-4 b	Yes	14.8	1	1.4 TeV	625 GeV < $m(\tilde{t}_1) < 850 \text{ GeV}$		
$\tilde{g}, \tilde{q}, \tilde{t}_1 \rightarrow q\tilde{g}\tilde{t}_1$	0	2 jets + 2 b	Yes	15.4	1	410 GeV	850-510 GeV	
$\tilde{g}, \tilde{q}, \tilde{t}_2 \rightarrow q\tilde{g}\tilde{t}_2$	2 e, μ, τ	Yes	20.3	1	0.4-1.0 TeV	$\text{BR}(\tilde{g} \rightarrow \nu\tilde{g}) > 20\%$		
Other	Scalar charm, $\tilde{c} \rightarrow c\tilde{c}^0$	0	2 e, τ	Yes	20.3	1	510 GeV	$m(\tilde{t}_1)=200 \text{ GeV}$

*Only a selection of the available mass limits on new states or phenomena is shown.

10⁻¹ 1 Mass scale [TeV]

QCD for Precise Hadron Collider Phenomenology

The Anatomy of Hadron-Hadron Collisions



- ▶ Hadron colliders are messy environments
- ▶ Access to high-energy interactions occurs in head on collisions and are described by so called partonic hard cross sections
- ▶ Radiation from incoming and outgoing partons always present
- ▶ Also soft physics related to the Underlying events, among other low-energy effects

Partonic Cross Section in Perturbation Theory

$$\hat{\sigma}(\alpha_s, \mu_F, \mu_R) = [\alpha_s(\mu_R)]^{n_\alpha} \left[\underbrace{\hat{\sigma}^{(0)}}_{\text{LO}} + \frac{\alpha_s}{2\pi} \underbrace{\hat{\sigma}^{(1)}}_{\text{NLO}}(\mu_F, \mu_R) + \left(\frac{\alpha_s}{2\pi}\right)^2 \underbrace{\hat{\sigma}^{(2)}}_{\text{NNLO}}(\mu_F, \mu_R) + \dots \right]$$

from L. Dixon

Problem: Leading-order, tree-level predictions only **qualitative**

due to **poor convergence**

of expansion in $\alpha_s(\mu)$

(setting $\mu_R = \mu_F = \mu$)

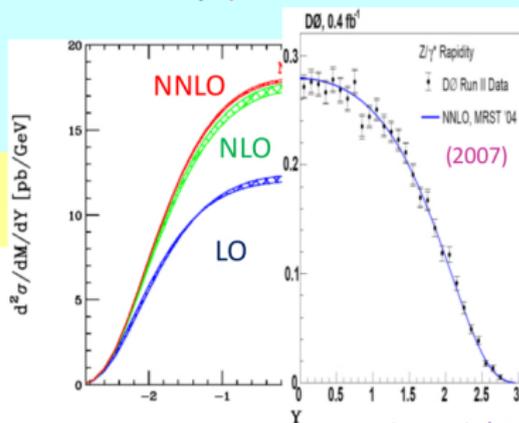
Example: Z production at Tevatron

Distribution in rapidity Y

$$Y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$$

$$\frac{d\sigma}{dY} \quad \text{has} \quad n_\alpha = 0$$

still ~50% corrections, LO \rightarrow NLO



[Anastasiou, Dixon, Melnikov, Petriello hep-ph/0312266]

Quantum Corrections in QCD

Goal:

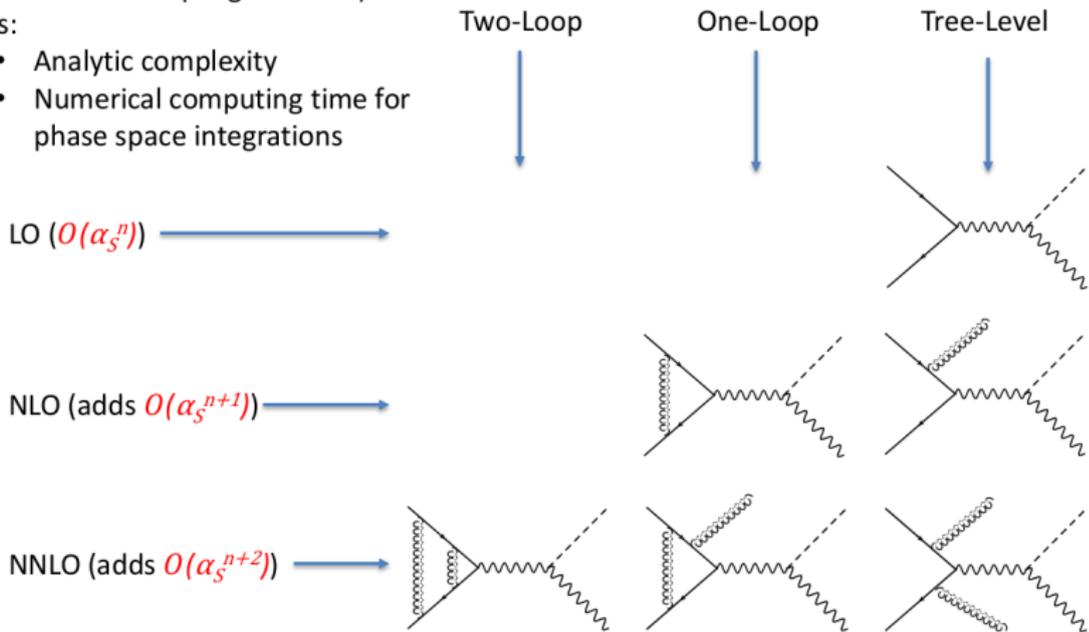
- Increased accuracy (expansion in small coupling constant)

Issues:

- Analytic complexity
- Numerical computing time for phase space integrations

$$d\sigma^{\text{NNLO}} = \alpha_S^n (m_{\text{tree}} + \alpha_S m_{1\text{-loop}} + \alpha_S^2 m_{2\text{-loop}})$$

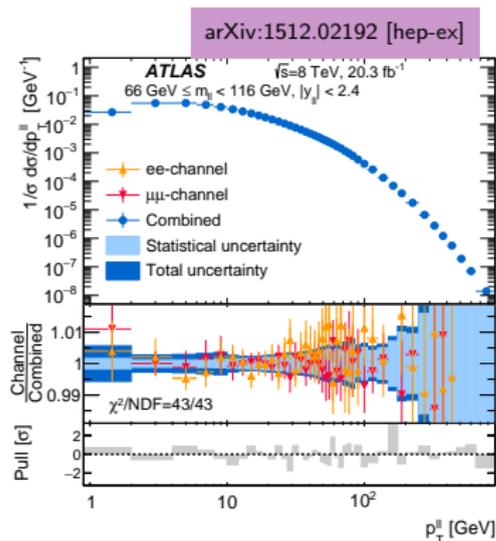
$$= \alpha_S^n (\tilde{m}_{\text{tree}} + \alpha_S \tilde{m}_{1\text{-loop}} + \alpha_S^2 \tilde{m}_{2\text{-loop}})$$



W. Giele

The $\sim 1\%$ Frontier at the LHC

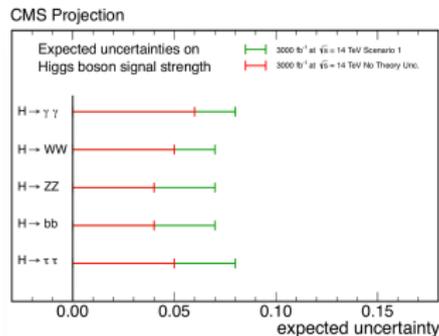
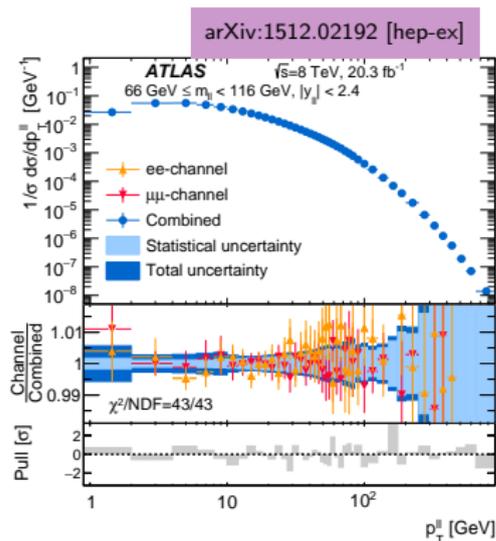
- ▶ p_T^{ll} in Drell-Yan, an impressive example of precise differential measurements by ATLAS
- ▶ By normalizing to inclusive Z cross section, improvement in uncertainties
- ▶ Total uncertainties below 1% for $p_T^{ll} < 200$ GeV



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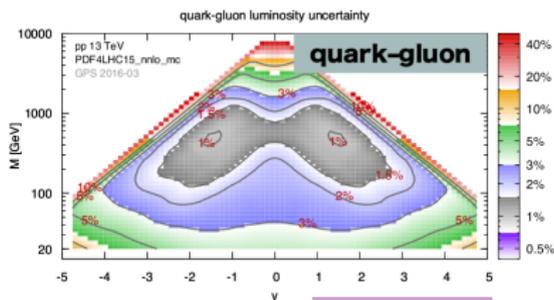
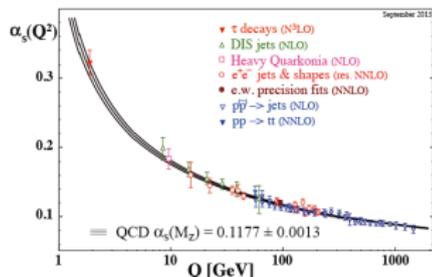
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- ▶ Higgs phenomenology will benefit also from high luminosity runs
- ▶ Signal strength relative errors could reach few percent
- ▶ Recent theoretical advances would allow matching this precision



Ingredients for QCD at $\sim 1\%$

In order to compute quantum QCD corrections two fundamental inputs are required: the strong coupling α_s and the *Parton Distribution Functions*



- ▶ Naively one is to expect NLO QCD corrections to be of order $\sim 10\%$ and NNLO QCD at $\sim 1\%$
- ▶ Perturbative calculations are also required for the partonic cross sections associated to the signal studied

State-of-the-Art QCD Phenomenology

dijets	$O(3\%)$	gluon-gluon, gluon-quark	PDFs, strong couplings, BSM
H+0 jet	$O(3-5 \%)$	fully inclusive (N3LO)	Higgs couplings
H+1 jet	$O(7\%)$	fully exclusive; Higgs decays, infinite mass tops	Higgs couplings, Higgs p_t , structure for the ggH vertex.
tT pair	$O(4\%)$	fully exclusive, stable tops	top cross section, mass, p_t , FB asymmetry, PDFs, BSM
single top	$O(1\%)$	fully exclusive, stable tops, t-channel	V_{tb} , width, PDFs
WBF	$O(1\%)$	exclusive, VBF cuts	Higgs couplings
W+j	$O(1\%)$	fully exclusive, decays	PDFs
Z+j	$O(1-3\%)$	decays, off-shell effects	PDFs
ZH	$O(3-5 \%)$	decays to bb at NLO	Higgs couplings (H-> bb)
ZZ	$O(4\%)$	fully exclusive	Trilinear gauge couplings, BSM
WW	$O(3\%)$	fully exclusive	Trilinear gauge couplings, BSM
top decay	$O(1-2 \%)$	exclusive	Top couplings
H -> bb	$O(1-2 \%)$	exclusive, massless	Higgs couplings, boosted

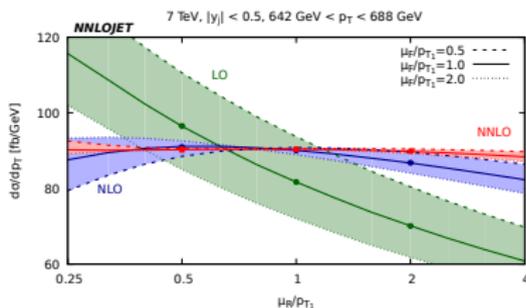
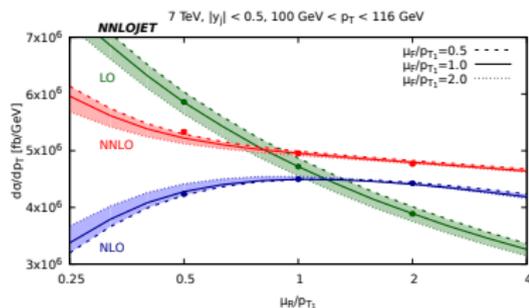
From K. Melnikov

Inclusive Jet Production @ NNLO QCD

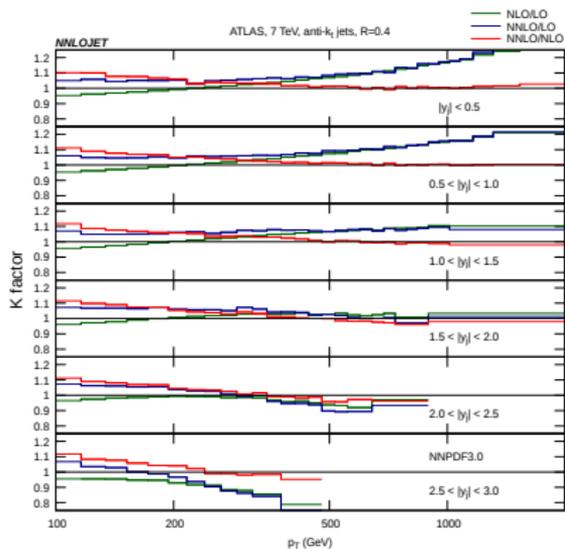
- ▶ Inclusive jet production is a fundamental process for hadron collider phenomenology
- ▶ It constrains directly the gluon PDFs, and in that way it has an impact on all theory predictions
- ▶ Although NNLO QCD PDFs appear in the market, employing data sets for inclusive jet production, approximations have been made as a full NNLO QCD calculation wasn't available
- ▶ Very recently [Currie](#), [Glover](#) and [Pires](#) ([arXiv:1611.01460](#)) have presented the first NNLO QCD results including all subprocesses
- ▶ This is the conclusion of a formidable task that started around 1999

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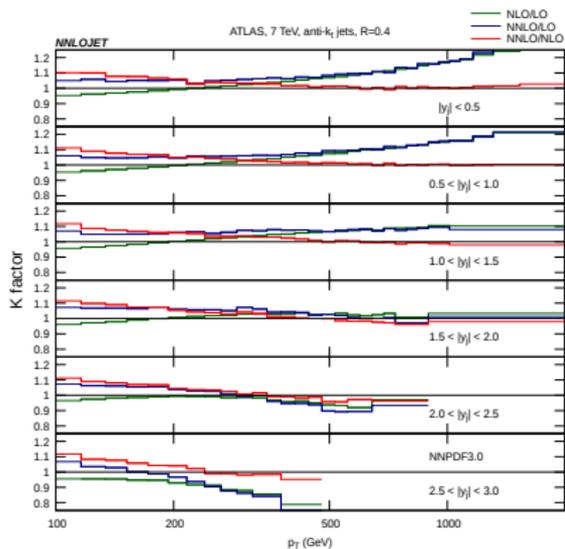


Structure of Corrections over PS

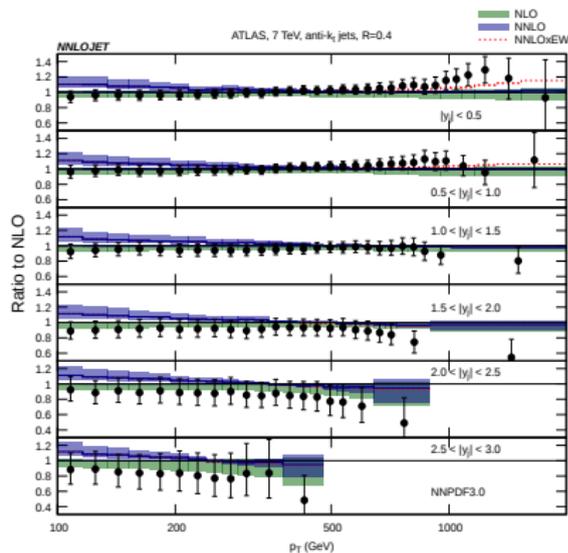


- ▶ Perturbative series converges well for large jet p_T
- ▶ But $\sim 10\%$ NNLO corrections around 100 GeV

Structure of Corrections over PS



- ▶ Perturbative series converges well for large jet p_T
- ▶ But $\sim 10\%$ NNLO corrections around 100 GeV



- ▶ Comparison to ATLAS 7 TeV data shows systematic deviations for low p_T
- ▶ We might see an impact on PDFs fits

Building NNLO QCD Corrections

- ▶ Two-loop amplitudes for process X
- ▶ One-loop amplitudes for process $X + g$
- ▶ Tree-level amplitudes for the processes $X + gg$, $X + q\bar{q}$, etc
- ▶ Strategy to handle and cancel IR divergences

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- ▶ Computing two-loop amplitudes is a significant challenge
 - ▶ Need fine control of numerical integration of one-loop and tree-level amplitudes over unresolved regions of PS
 - ▶ Procedures for extracting IR divergences (subtraction/slicing) can be cumbersome. A lot of recent progress: *antenna subtraction*, *q_T subtraction*, *N -jettiness slicing*, *sector-improved residue subtraction*, among other

High Multiplicity Amplitude Calculations

Numerical Unitarity for Computing Amplitudes

AIM: Write amplitude (\mathcal{A}) as a sum of master integrals.

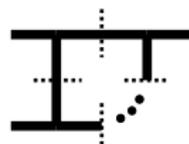
$$\mathcal{A} = \int A = \int \sum_i \frac{N_i}{\rho^1 \dots \rho^{n_i}} = \sum_i c_i \int \frac{t_i^{\text{master}}}{\rho^1 \dots \rho^{n_i}}$$

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General algorithm:



Unitarity \longleftrightarrow Residue (A) $\xleftrightarrow{\text{Subtraction}}$ $N_i = \sum_i c_i t_i$

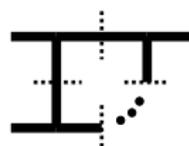
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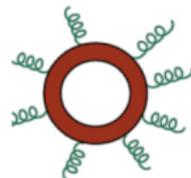
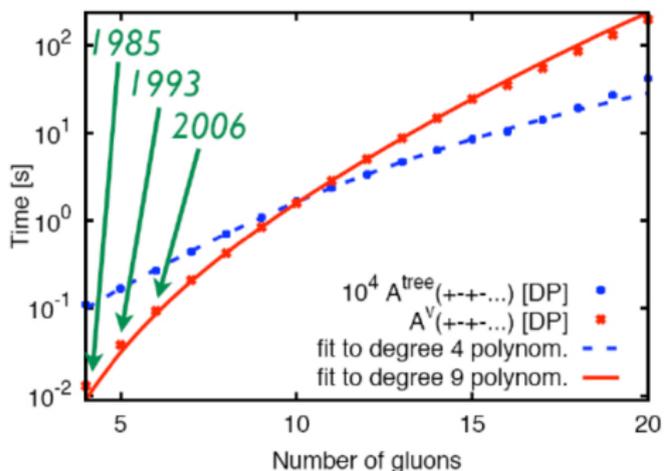
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2-loop **complications**:

- ▶ IBPs - how to find basis $\{t_i^{\text{master}}, t_i^{\text{surface}}\}$? [Ita 15]
- ▶ Much richer structure of cuts and master integrals
- ▶ Handle efficiently the regressions of tensor coefficients

For 1-loop Amplitudes, A Powerful Technique!



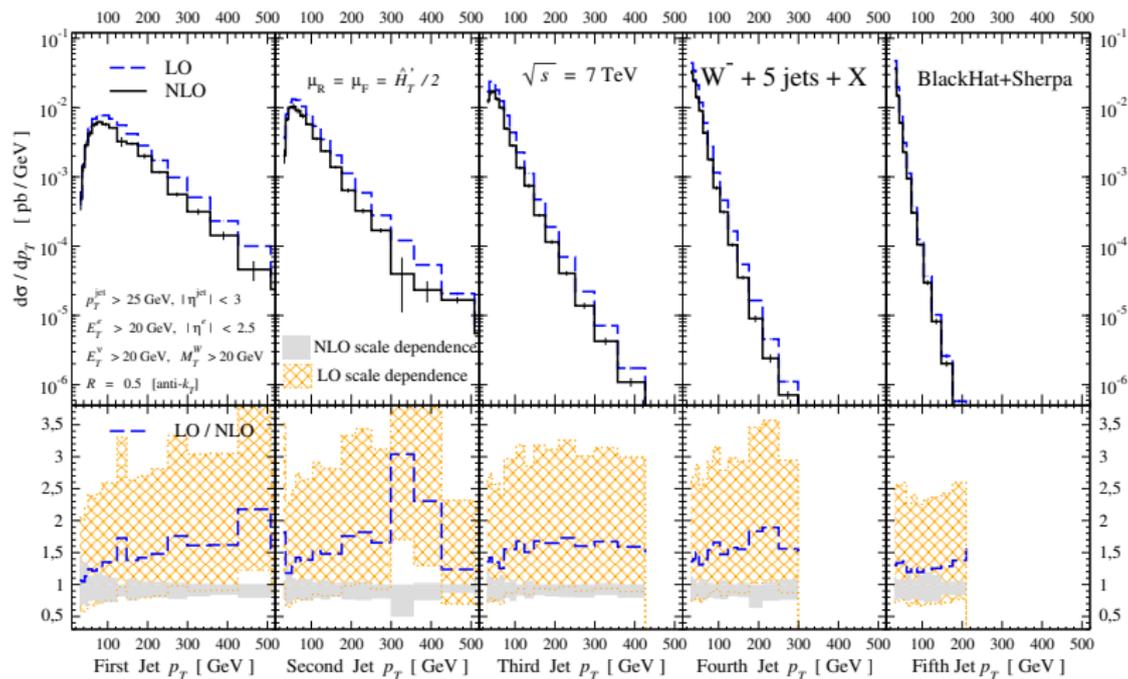
[Giele, Zanderighi
arXiv:0806.2152]

BUT STILL VERY COMPUTER INTENSIVE

[BlackHat + Sherpa]

NTUPLES: STORE AS MUCH INFO AS POSSIBLE DURING YOUR COMPUTATION!

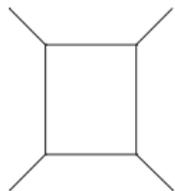
Jet p_T Spectra at NLO for $W + 5$ -Jet Production



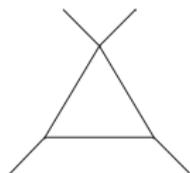
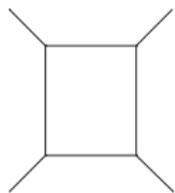
- ▶ Involves 1-loop amplitudes with 8 particles attached to the loop
- ▶ Real radiation with integration over PS of 6(7) particles

- ▶ Impressive improvement on the perturbative prediction
- ▶ Allows for tests of QCD in highly complex kinematic configurations

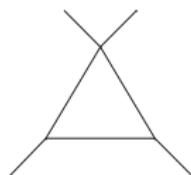
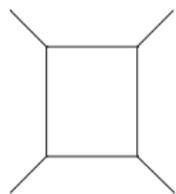
Cuts at One- and Two-Loops in 4-pt Amplitudes



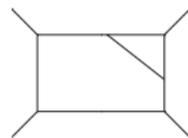
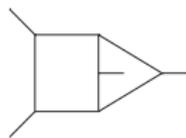
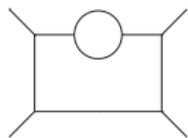
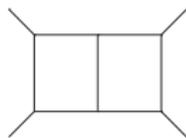
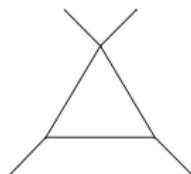
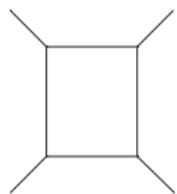
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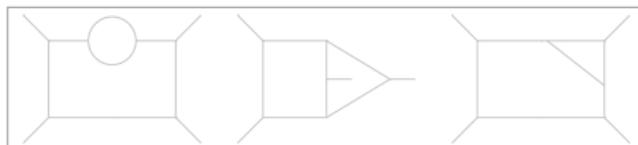
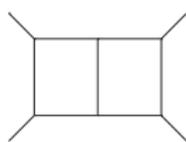
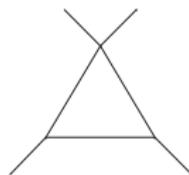
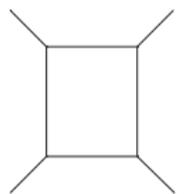
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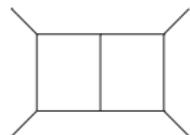
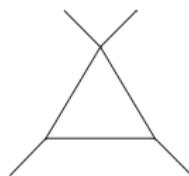
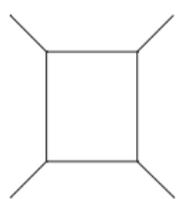
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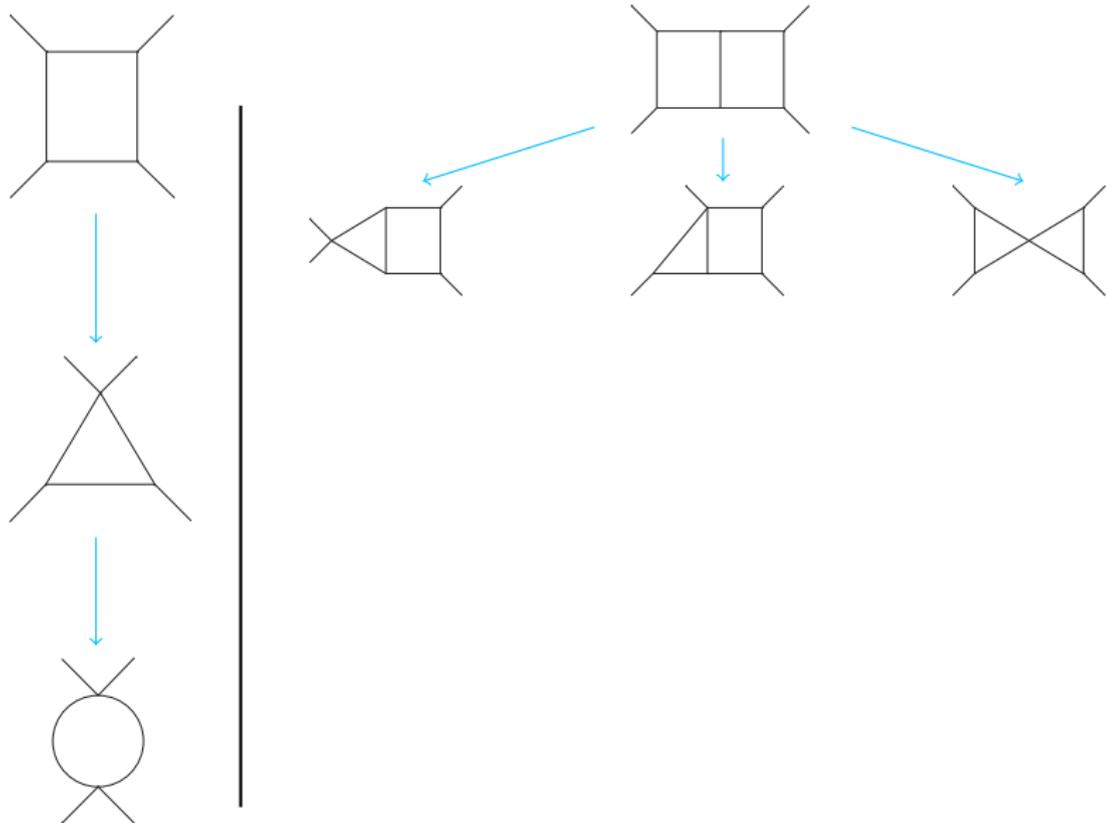
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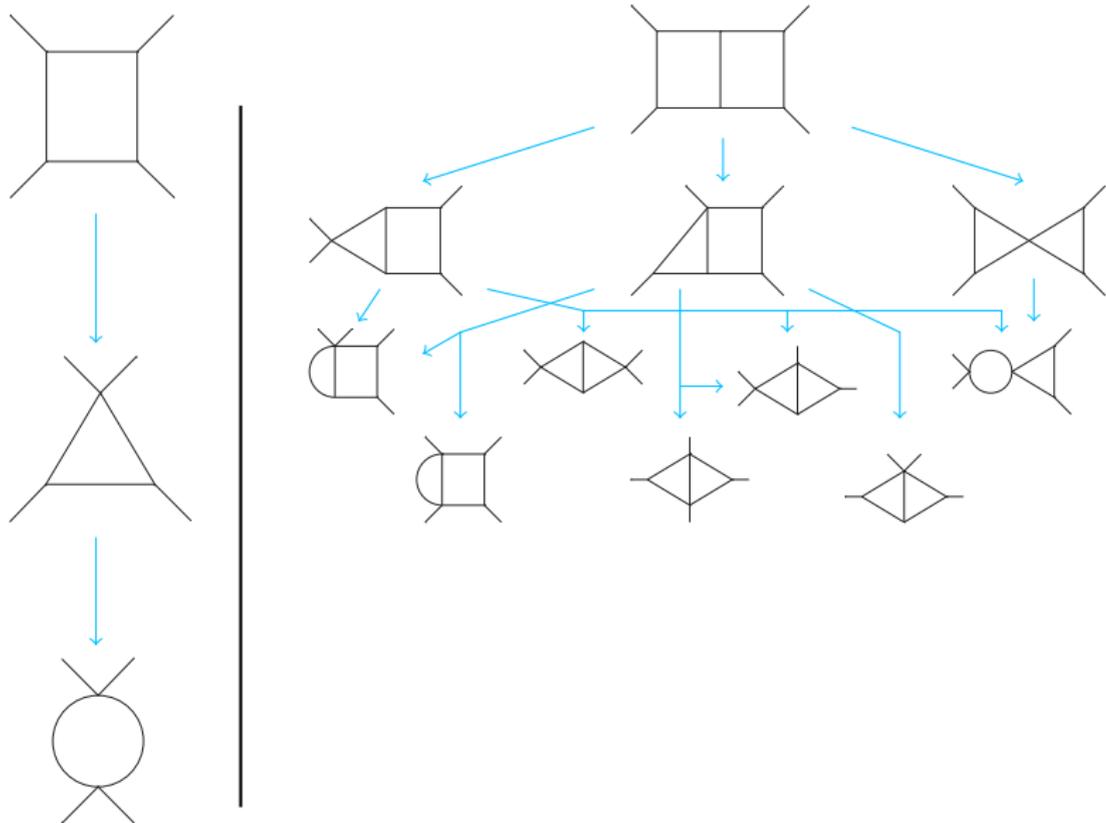
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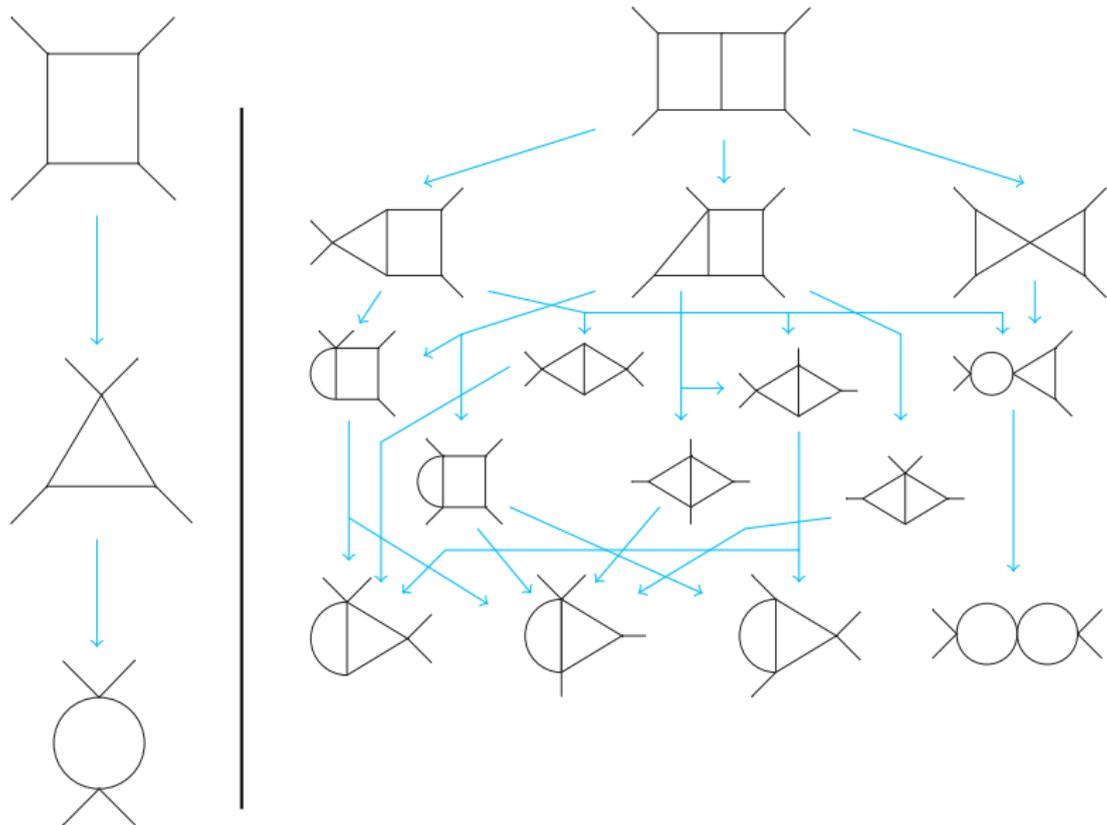
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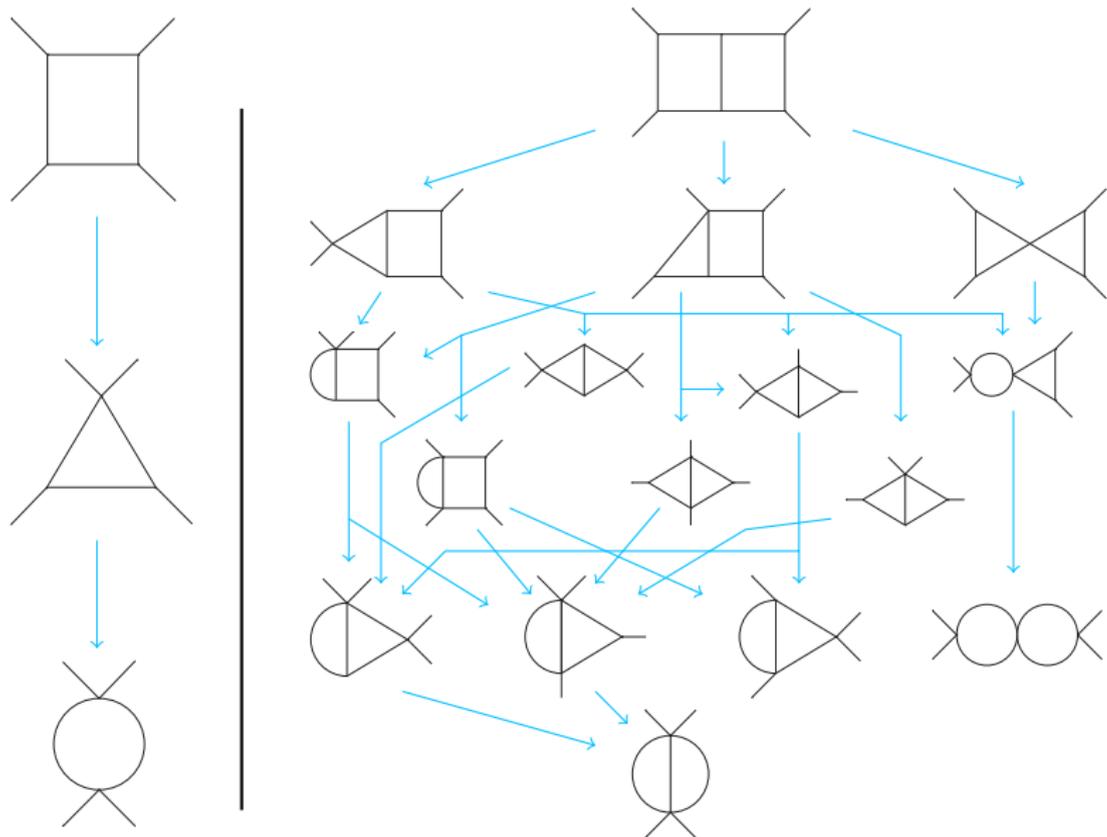
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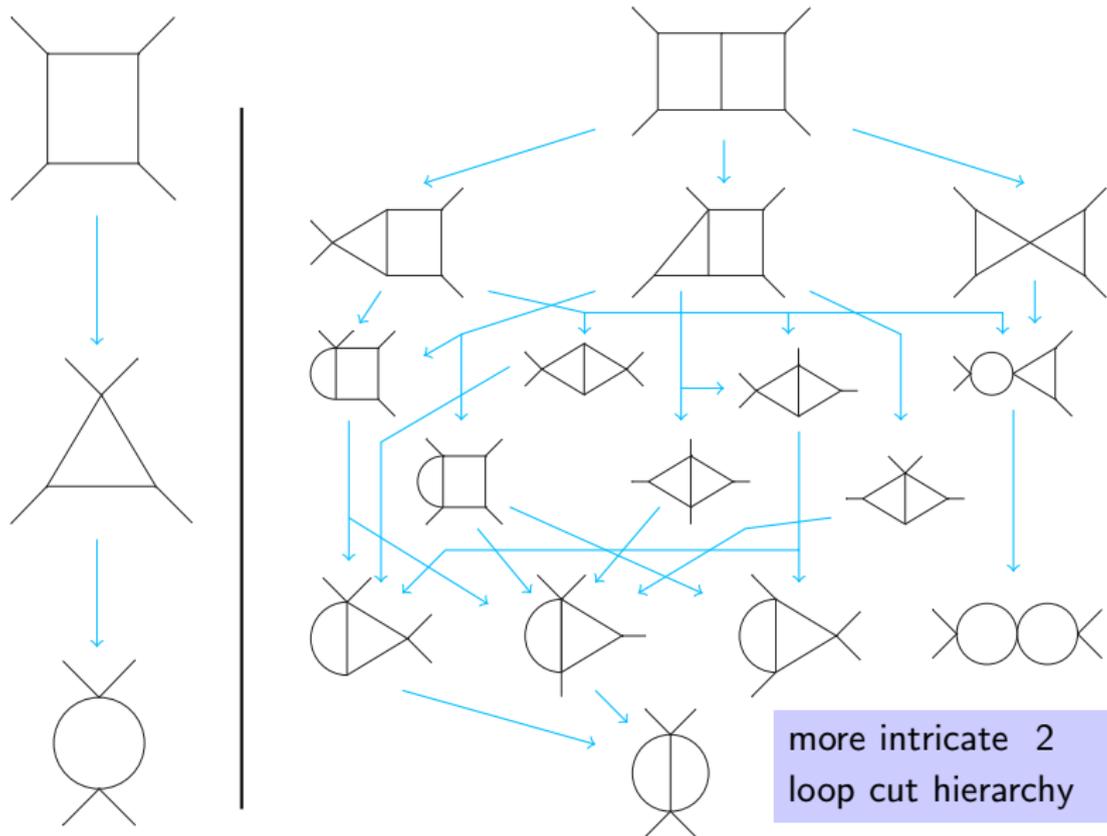
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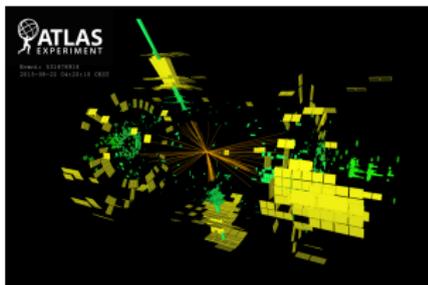


The Path to 5-pt Two-Loop Amplitudes

- ▶ First examples of amplitudes have started to appear
[Badger, Frellesvig, Mogull, Ochirov, O'Connell, Zhang], [Gehrmann, Henn, Lo Presti], [Dunbar, Jehu, Perkins]
- ▶ Important progress on integrand decomposition
[Ita], [Zhang, Larsen], [Mastrolia, Peraro, Primo, Bobadilla]
- ▶ 5-pt (master) integrals also appearing
[Papadopoulos, Tommasini, Wever], [Gehrmann, Henn, Lo Presti]

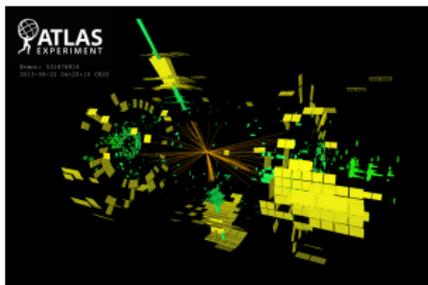
Conclusions

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Thanks!